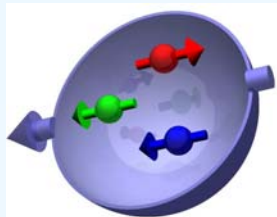


クォーク・グルーオンレベルで見た 核子スピンの起源

岩田高広
山形大学 理学部

核子スピンの何が起こったか？

クォークモデル

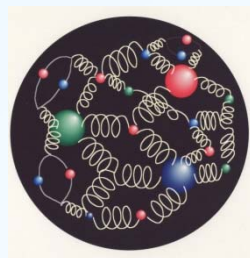


核子スピン：クォークスピンの合成

クォークスピン寄与は100%

QCDパートン描像

valence quarks +
sea-quarks
& gluons



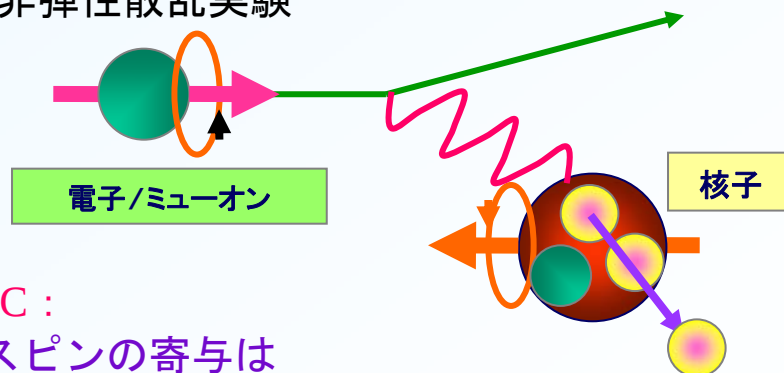
クォークスピン
30%



$Q^2=3\text{GeV}^2$

残りの寄与
70%

偏極深部非弾性散乱実験



1988 EMC :
クォークスピンの寄与は
非常に少ない! → spin crisis

1990-2008 SMC, SLAC, HERMES, COMPASS :
クォークスピン寄与は約30% (@ $Q^2=3\text{GeV}^2$)

Missing Spin の問題

(spin puzzle / Dark Spin)

核子スピンを担う候補

核子のスピン和則

$$\Delta\Sigma = \sum_i \Delta q_i^{\uparrow} - \Delta q_i^{\downarrow} + \Delta \bar{q}_i^{\uparrow} - \Delta \bar{q}_i^{\downarrow}$$

$$\frac{1}{2} = \frac{1}{2} \Delta\Sigma_{\text{quark spin}} + \Delta G_{\text{gluon spin}} + \langle L_Z \rangle_{\text{orbital}}$$

Missing Spinの候補

$$\Delta G = G^{\uparrow} - G^{\downarrow}$$

1) グルーオンスピン

グルーオン : QCDにおけるゲージ粒子(スピン1)

核子の4元運動量の半分近くを運ぶ

核子中でのグルーオン偏極度は？

2) クォーク・グルーオンの軌道回転

クォークモデルではクォークの軌道回転は無い

QCDパートン描像では、どうなるか？



CERNでの11カ国の国際共同研究

日本グループ：山形大学、宮崎大学、KEK、中部大学

偏極深部非弾性散乱実験

160GeV 偏極ミュオンビーム (SPS) & 偏極核子標的

スペクトロメータ：散乱ミュオン + ハドロンの同時検出

主な研究課題

- 1) グルーオン偏極度の測定
直接測定 / 間接測定
- 2) クォークスピン寄与のフレーバー分解、Valence/sea 分解
- 3) クォークの横運動量依存分布 (Sivers関数)
- 4) クォークの横方向偏極効果 (transversity、Collins効果)
- 5) Λ 偏極

COMPASS SETUP



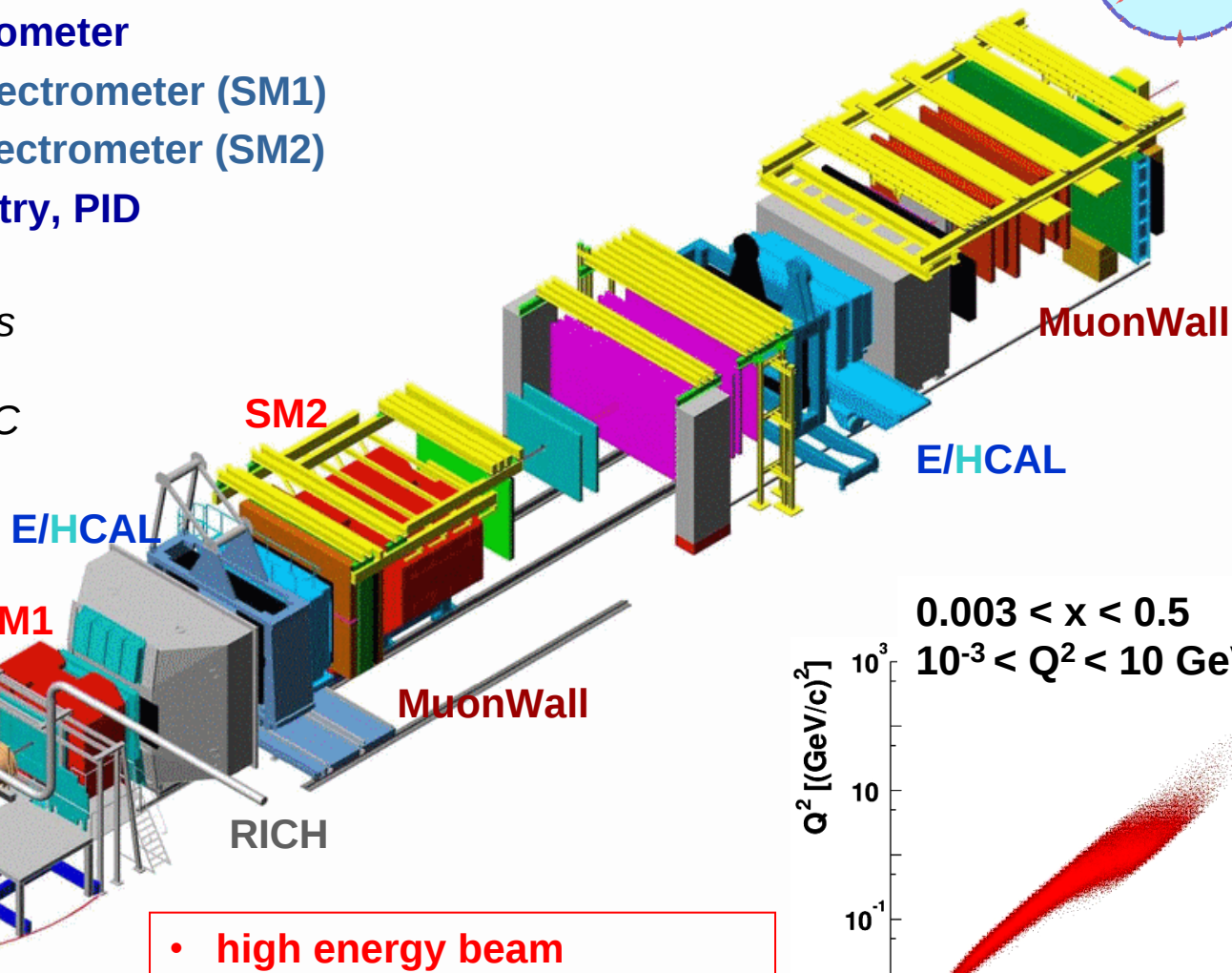
two stages spectrometer

- Large Angle Spectrometer (SM1)
- Small Angle Spectrometer (SM2)

tracking, calorimetry, PID

SciFi
Silicon
Micromegas
GEMs

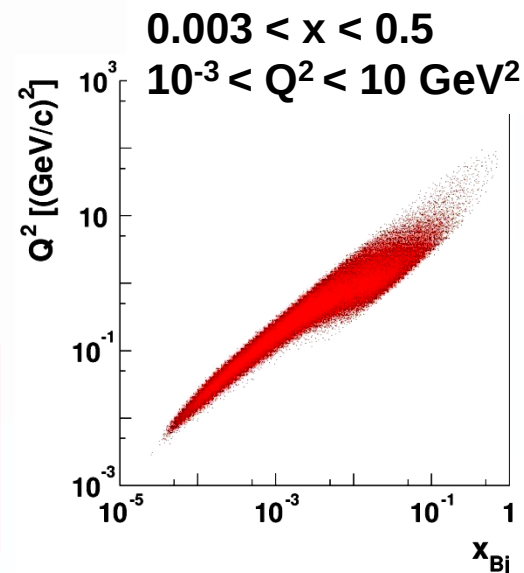
Straws
SDC
MWPC
W45



Pol.
 μ beam

$E_\mu = 160 \text{ GeV}$
 $I_\mu = 2 \times 10^8 / \text{spill}$,
 spill = 4.8s, cycle = 16.8s
 $P_\mu = 80\%$

- high energy beam
- broad kinematical range
- large angular acceptance

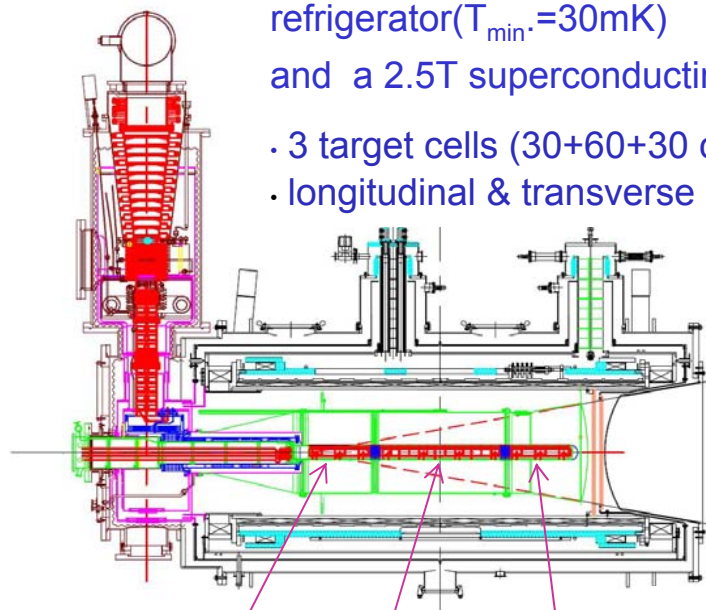


THE POLARIZED TARGET SYSTEM

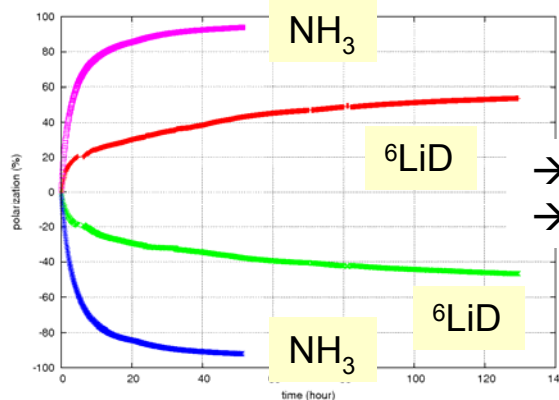
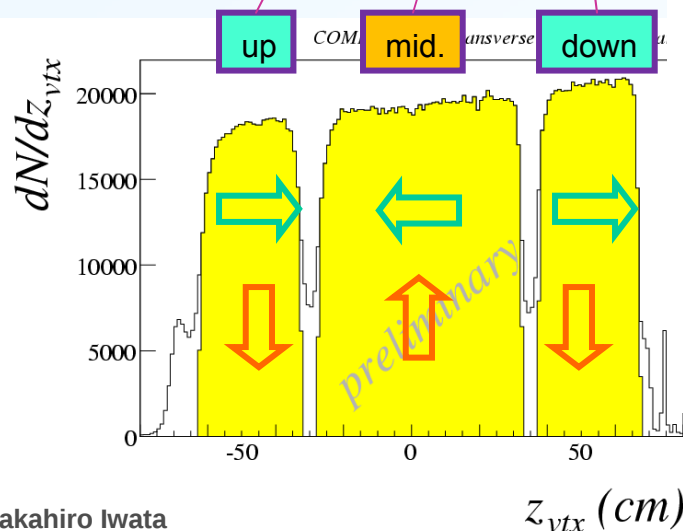


DNP technique with a dilution refrigerator ($T_{\min.}=30\text{mK}$) and a 2.5T superconducting magnet

- 3 target cells (30+60+30 cm)
- longitudinal & transverse modes



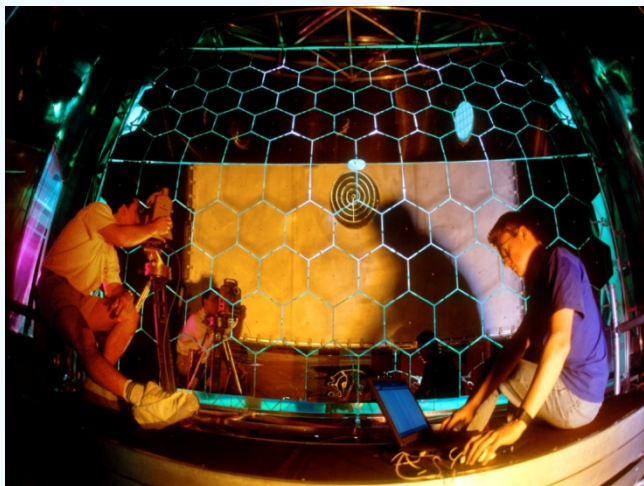
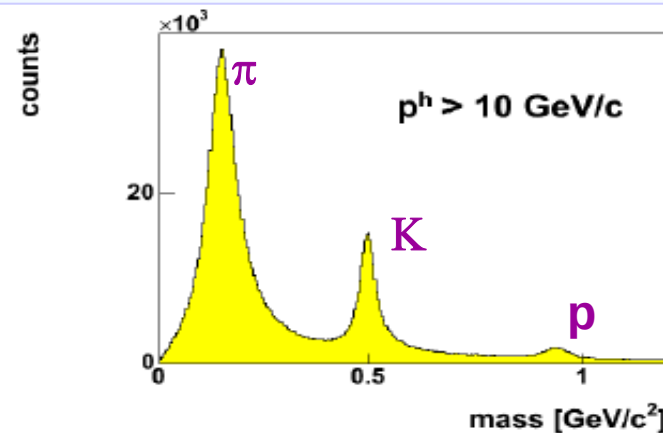
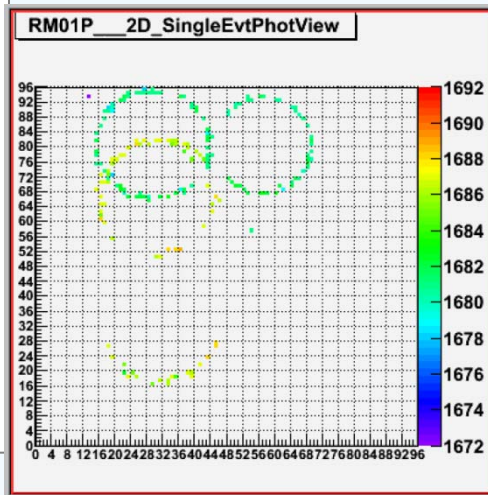
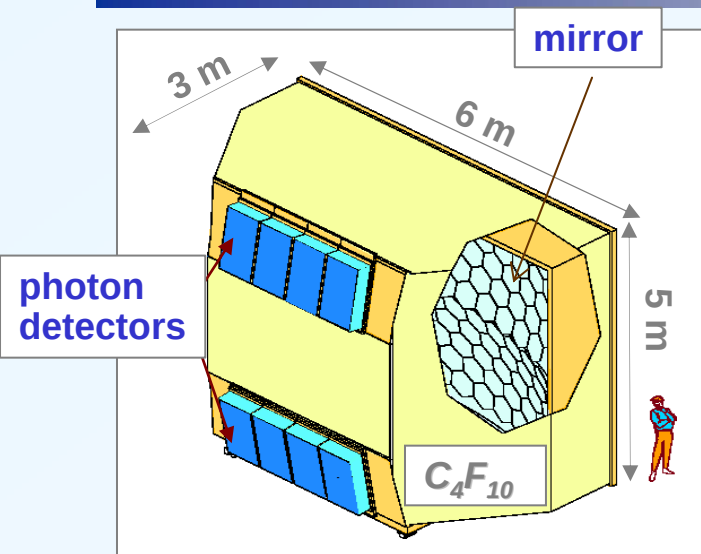
世界最大の
偏極標的



- Polarization of NH₃ ~90%
- Polarization of ⁶LiD ~50%

• ⁶Li (~ α +d) also polarized
→ dilution factor f=50%

THE RICH DETECTOR



- radiator gas: C_4F_{10}
- mirror wall: 20 m² surface
- photon-detectors:
 - outer part (75%) MWPC(pad RO) with CsI cathode
 - inner part(25%) 576 MAPMTs with indiv. telescope

threshold momenta

- $p_\pi = 2 \text{ GeV/c}$
- $p_K = 9 \text{ GeV/c}$
- $p_p = 17 \text{ GeV/c}$

Installed in 2005,
Used in data taking from 2006

DIRECT MEASUREMENTS OF $\Delta G/G$



グルーオン偏極度: $\frac{\Delta G}{G}$

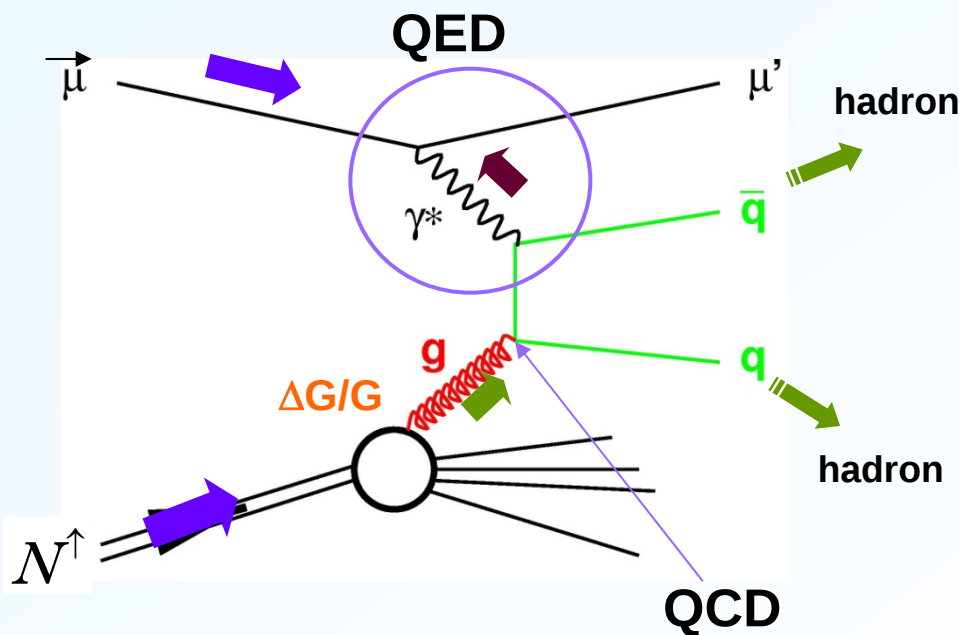
直接測定

グルーオンが関与する反応の非対称度

Theory (QCD)

$$A_{LL} = \langle \hat{a} \rangle \times \frac{\Delta G}{G}$$

Photon Gluon Fusion: $\gamma + g \rightarrow q \bar{q}$



$q = u, d, s$

“HIGH p_T HADRON PAIRS”

2 hadrons with high p_T

😊 large statistics

☹ physical BG: MC dependent

• hard scale: p_T

$q = c$

“OPEN CHARM”

one charm hadron detected

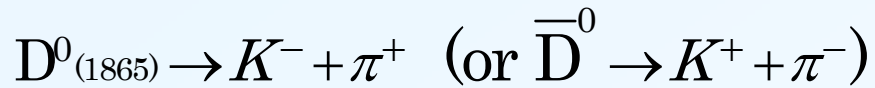
😊 small phys. BG, less MC dependent

☹ small statistics

• hard scale: charm mass, $\mu^2 = 4m_c^2$

• COMPASS original

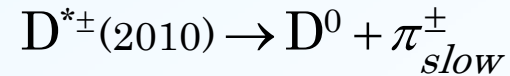
$\Delta G/G$ FROM OPEN CHARM



D^* tagging

BR:4%

Kaon id. by RICH

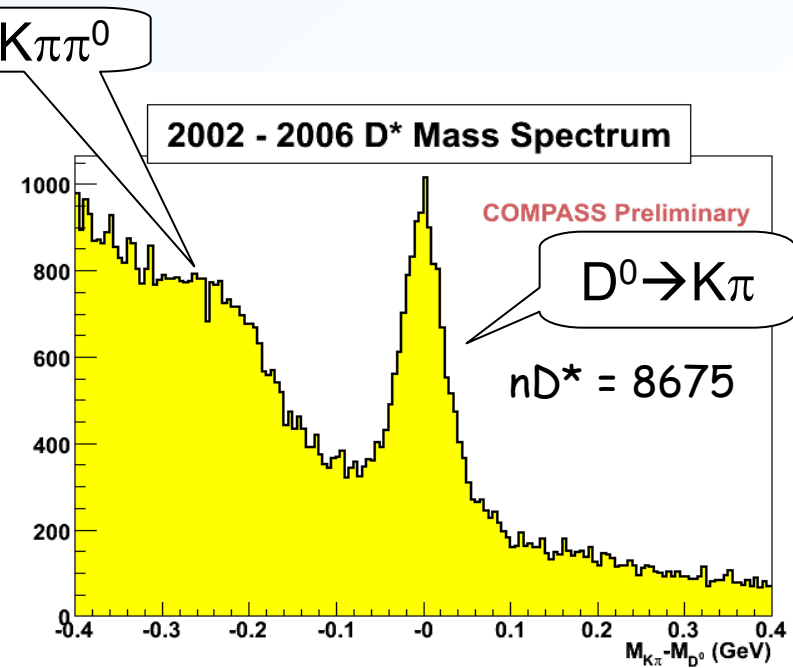
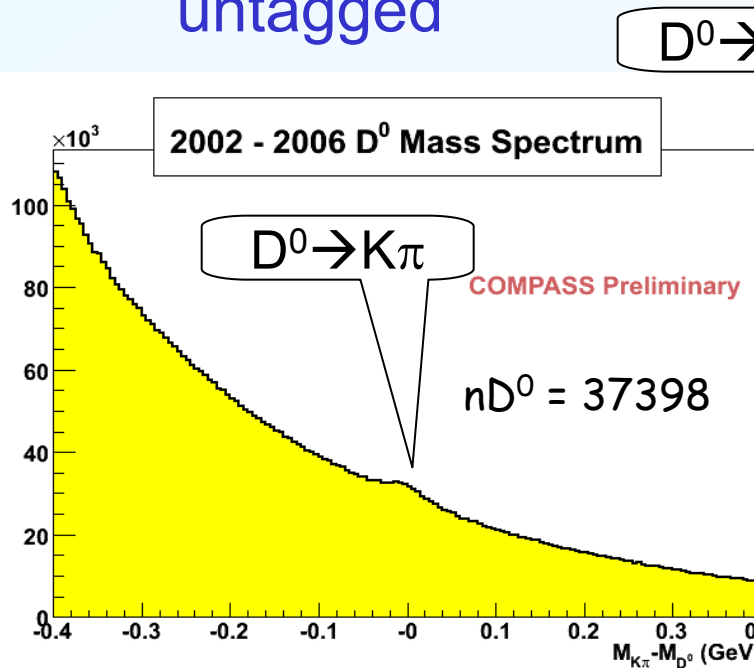


BR:68%



$K + \pi$

untagged



$M_{K\pi} - M_{D^0} \text{ (GeV)}$

$M_{K\pi} - M_{D^0} \text{ (GeV)}$

INDIRECT MEASUREMENT OF $\Delta G/G$ COMPASS



- Polarized inclusive DIS

$$e^\uparrow + N^\uparrow \rightarrow e' + X$$

間接測定

→ polarized structure function: $g_1(x, Q^2)$

at LO-QCD $g_1(x, Q^2) = \frac{1}{2} \sum_i e_i^2 \Delta q_i$ ← クォーク偏極

at NLO-QCD

$$g_1(x, Q^2) = \frac{1}{9} \left(C_{NS} \otimes \left[\pm \frac{3}{4} \Delta q_3 + \frac{1}{4} \Delta q_8 \right] + C_S \otimes \Delta \Sigma + C_G \otimes \Delta G \right)$$

クォーク偏極

グルーオン偏極

$$\Delta q_3 = \Delta u - \Delta d$$

$$\Delta q_8 = \Delta u + \Delta d - 2\Delta s$$

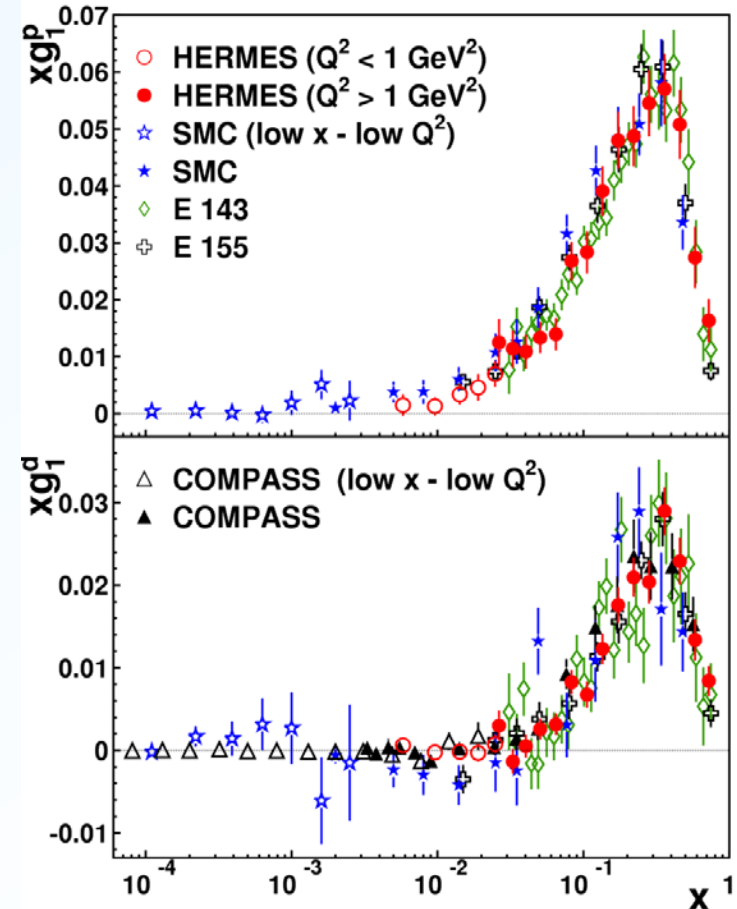
C_{NS}, C_S, C_G are Wilson coefficients given by QCD

Fit world data with QCD DGLAP eq. ,

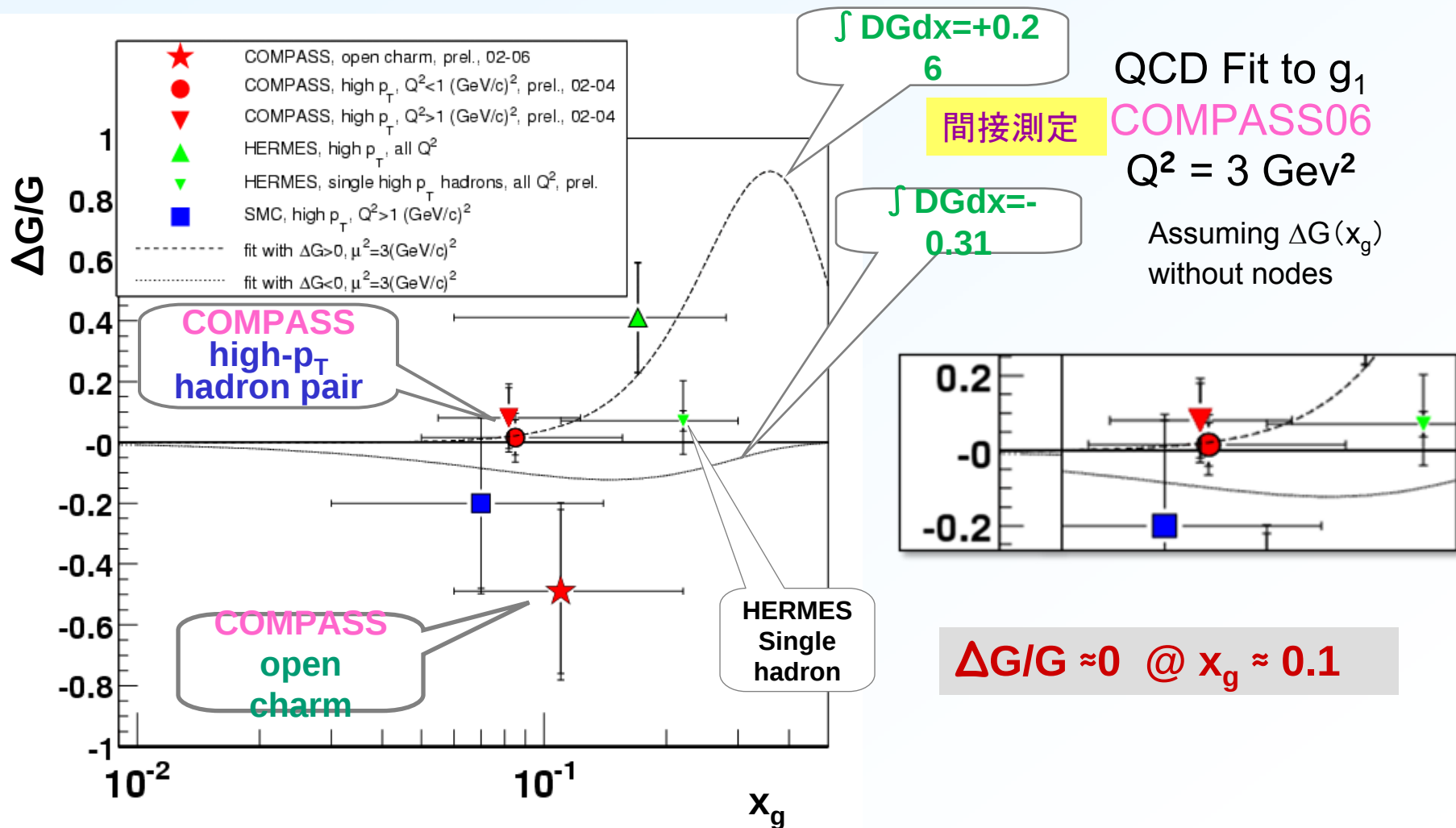
$$\Delta \Sigma(x) = \sum_i \Delta q_i : \text{クォーク偏極の和}$$

$$\Delta G(x) : \text{グルーオン偏極}$$

が同時に得られる



RESULTS FOR $\Delta G/G$



RECALL: charm point at larger scale ($\mu^2 = 13 \text{ GeV}^2$)

CONTRIBUTIONS TO NUCLEON SPIN

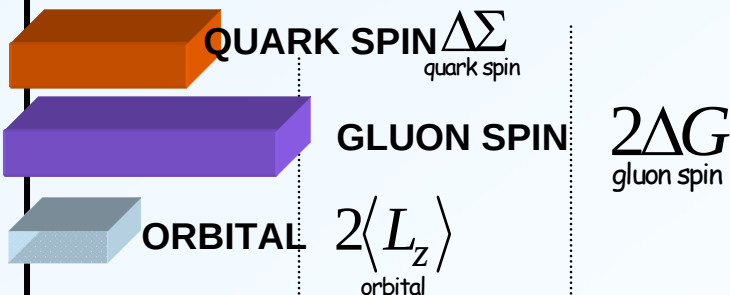


$$\frac{1}{2} = \frac{1}{2} \Delta\Sigma_{\text{quark spin}} + \Delta G_{\text{gluon spin}} + \langle L_z \rangle_{\text{orbital}} \rightarrow 100\% = \Delta\Sigma_{\text{quark spin}} + 2\Delta G_{\text{gluon spin}} + 2\langle L_z \rangle_{\text{orbital}}$$

e.g. COMPASS 06

$\Delta G > 0$ solution

evaluated



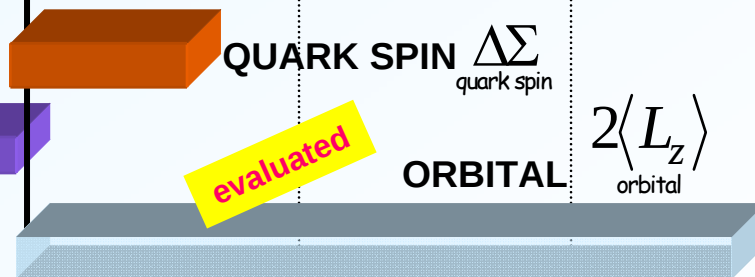
核子スピンの半分は
グルーオンスピンの寄与

$\Delta G < 0$ solution

$2\Delta G_{\text{gluon spin}}$

GLUON SPIN

evaluated



負のグルーオンスピ
ン寄与を打ち消すた
めに大きな軌道回転
寄与が必要

-50%

+50%

+100%

+150%

Negative contribution

Positive contribution

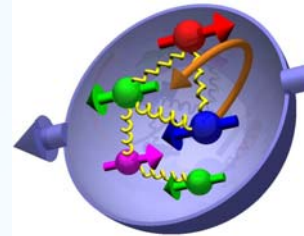
- ・グルーオンスピンの寄与は本質的
- ・軌道回転寄与も無視できない

COMPASSでの今後の研究の方向性



- グルーオンスピン寄与の確定
 - 蓄積データの解析、 新たなデータ収集(統計精度の向上)

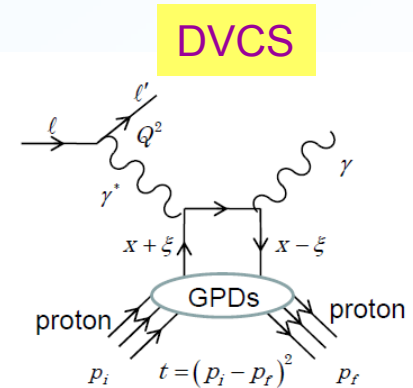
- **クォークの軌道回転寄与の研究**



- **一般化パートン分布(GPD)**

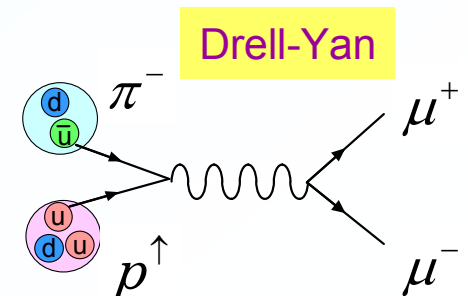
- 深部仮想コンプトン散乱(DVCS)の検出
- **Jiの全角運動量和則**からクォークの軌道回転寄与を抽出

$$J^q = \frac{1}{2} \Delta q + L_z = \lim_{t \rightarrow \infty, \xi \rightarrow 0} \frac{1}{2} \int_{-1}^1 x \left[H^q(x, \xi, t) + E^q(x, \xi, t) \right] dx$$



- **偏極ドレリヤン**

- Boer-Mulder分布関数(クォークのL・S結合)
- Sivers分布関数(クォークの軌道回転効果)

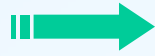


BACKUP SLIDES

QCD FIT WITH g_1 DETAILS

Start from PDFs at initial $Q^2=Q_0^2$

$\Delta\Sigma, \Delta q_3, \Delta q_8, \Delta G$



$$\Delta q_k(x) = \eta_k \frac{x^{\alpha_k} (1-x)^{\beta_k} (1+\gamma_k x)}{\int_0^1 x^{\alpha_k} (1-x)^{\beta_k} (1+\gamma_k x) dx}$$

constraints

$$\eta_3 = F + D$$

$$\eta_8 = 3F - D$$

$$\gamma_3 = \gamma_8 = \gamma_G = 0$$

DGLAP: to go from Q_0^2 to Q_{data}^2

non-singlet

$$\frac{d}{dt} \Delta q_{NS} = \frac{\alpha_s(t)}{2\pi} P_{qq}^{NS} \otimes \Delta q_{NS}$$

$$\Delta q_{NS} = (3/4) [\Delta q_3 + (1/3) \Delta q_8]$$

singlet & gluon

$$\frac{d}{dt} \begin{bmatrix} \Delta\Sigma \\ \Delta G \end{bmatrix} = \frac{\alpha_s(t)}{2\pi} \begin{bmatrix} P_{qq}^S & 2n_f P_{qG}^S \\ P_{Gq}^S & P_{GG}^S \end{bmatrix} \otimes \begin{bmatrix} \Delta\Sigma \\ \Delta G \end{bmatrix}$$

where $t = \log(Q^2/\Lambda^2)$ and P_{qq}, P_{qG}, P_{Gq} are polarized splitting functions

g_1^{fit} at the Q_{data}^2

at NLO

$$g_1^{fit}(x, Q^2) = \frac{1}{9} \left(C_{NS} \otimes \left[\pm \frac{3}{4} \Delta q_3 + \frac{1}{4} \Delta q_8 \right] + C_S \otimes \Delta\Sigma + C_G \otimes \Delta G \right)$$

where C_{NS}, C_S, C_G are Wilson coefficients

minimize

$$\chi^2 = \sum_i \frac{[g_1^{fit}(x_i, Q^2) - g_1^{exp}(x_i, Q^2)]^2}{[\sigma(x_i, Q^2)]^2}$$

get

$(\alpha_k, \beta_k, \gamma_k, \eta_k)$

V.Yu.Alexakhin et al., Phys.Lett.B647(2007)8-17

QUARK DISTRIBUTION STUDY IN DY

unpol. DY

$$\hat{k} \propto \bar{h}_{1q}^{\perp(1)}(x_1) \cdot h_{1q}^{\perp(1)}(x_2)$$

FF free

cos2 ϕ modulation relates to Boer-Mulders function

single polarized DY

B.M. function

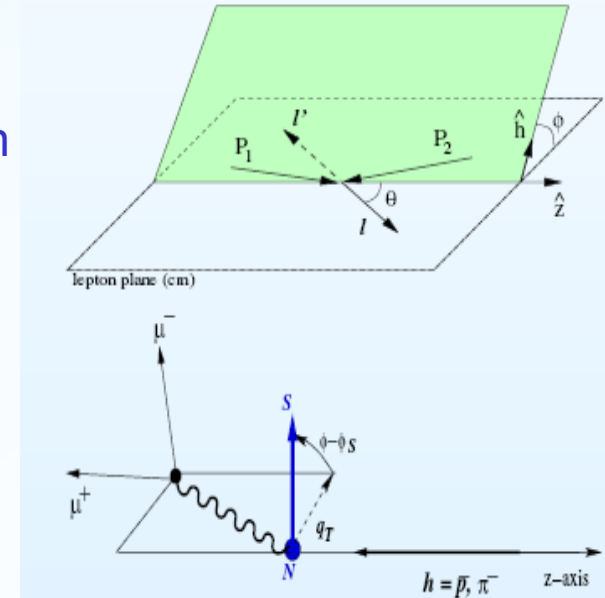
Transversity

$$A^{\sin(\phi+\phi_S)} \propto \bar{h}_{1q}^{\perp(1)}(x_1) \cdot h_{1q}(x_2)$$

$$A^{\sin(\phi-\phi_S)} \propto f_1^{\bar{q}}(x_1) \cdot f_{1T}^{\perp(1)q}(x_2)$$

Unpol. PDF

Sivers function



prediction

$$f_{1T}^{\perp} |_{DY} = -f_{1T}^{\perp} |_{SIDIS}$$

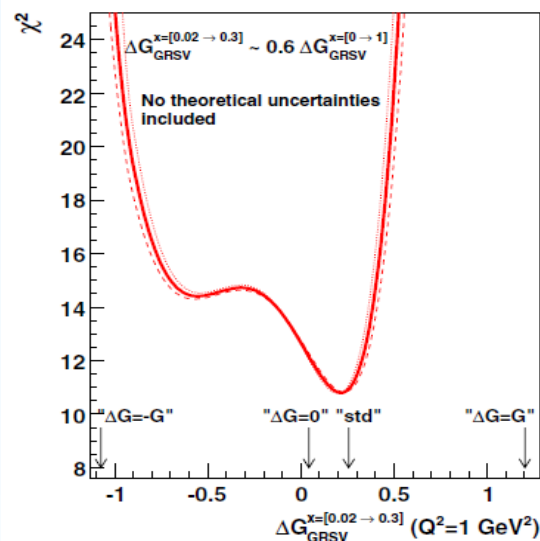
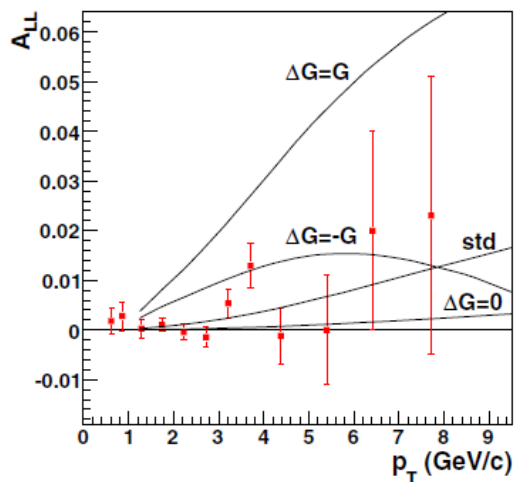
Collins , PL **B536** (2002) 43

$\Delta G/G$ FROM RHIC-SPIN

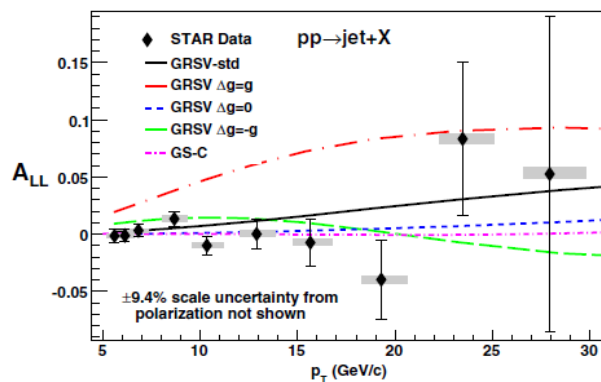


PHYSICAL REVIEW D 76, 051106(R) (2007)

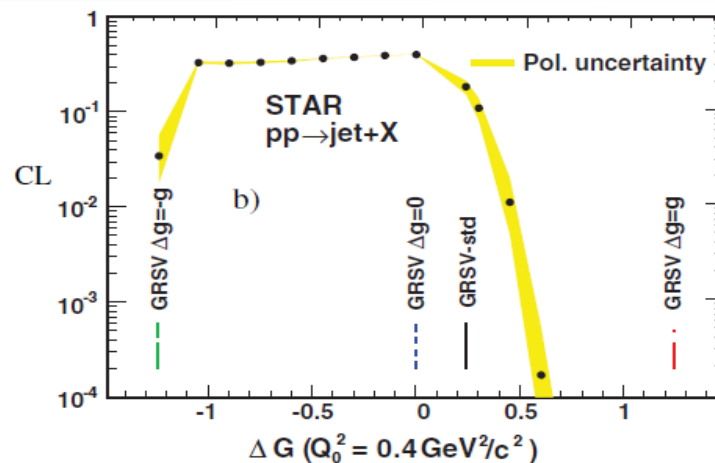
$$p^\uparrow + p^\uparrow \rightarrow \pi^0 + X$$



$$p^\uparrow + p^\uparrow \rightarrow jet + X$$



PRL 100, 232003 (2008)



TRANSVERSE MOMENTUM DEPENDENT PDFs

Nucleon

unpol.

long. pol.

trans. pol.

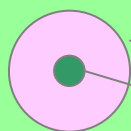
unpol.

long. pol.

trans. pol.

$$f_1 (= q)$$

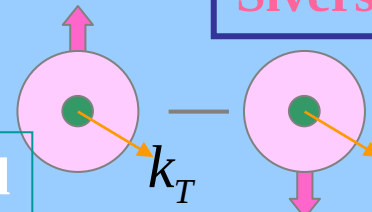
number
density



Nucleon
Parton

$$f_{1T}^\perp$$

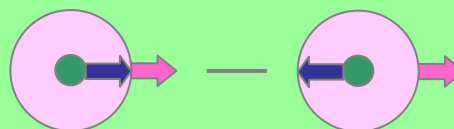
Sivers



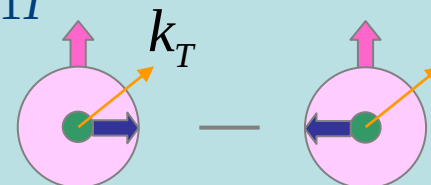
T-odd

$$g_{1L} (= \Delta q)$$

helicity



$$g_{1T}$$



$$h_1 (= \Delta_T q)$$

transversity



$$h_1^\perp$$

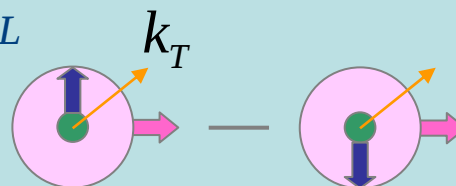
**Boer-
Mulders**



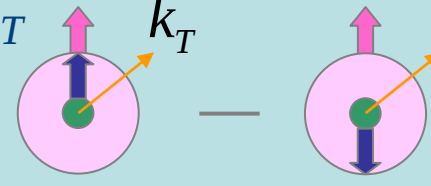
T-odd

survive after
integration over k_T

$$h_{1L}^\perp$$

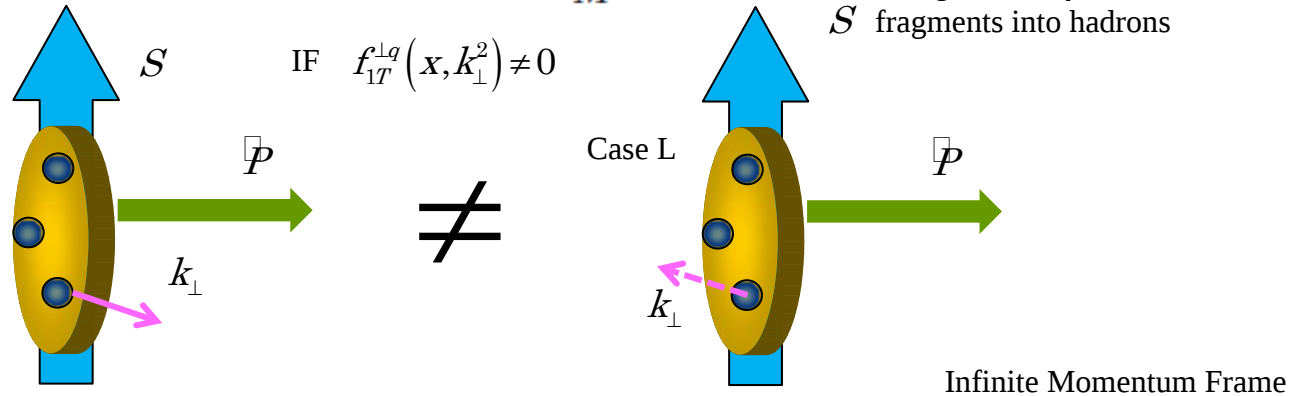


$$h_{1T}^\perp$$



Sivers: distribution of unpol. quarks in $\perp$ pol. proton

$$f_{q/p^\uparrow}(x, \mathbf{k}_\perp) = f_1^q(x, \mathbf{k}_\perp^2) - f_{1T}^{\perp q}(x, \mathbf{k}_\perp^2) \frac{(\hat{\mathbf{P}} \times \mathbf{k}_\perp) \cdot \mathbf{S}}{M}$$



Boer-Mulders: distribution of $\perp$ pol. quarks in unpol. proton

$$f_{q^\uparrow/p}(x, \mathbf{k}_\perp) = \frac{1}{2} \left[f_1^q(x, \mathbf{k}_\perp^2) - h_1^{\perp q}(x, \mathbf{k}_\perp^2) \frac{(\hat{\mathbf{P}} \times \mathbf{k}_\perp) \cdot \mathbf{S}_q}{M} \right]$$

